

# Invariant Solution of the no Stationary Elastic Deformation of Hollow Cylinder Problem

R.P. Abdrahmanova  
 Department of Computer Science and Robotics  
 Ufa State Aviation Technical University  
 Ufa, Russia  
 e-mail: kickufa@online.ru

## Abstract<sup>1</sup>

Consider the elementary case of the no stationary elastic deformation of hollow cylinder problem under action of internal and external pressure. In exact statement is gotten the problem with free boundaries. Using group of the transformations admitted by the equation some decisions of a problem dependent on infinite number of constant are found.

## 1. Introduction

Consider the equation

$$\frac{\partial^2 v}{\partial t^2} = \frac{\partial}{\partial r} \left( \frac{\partial v}{\partial r} + r^{-1} v \right) \quad (1)$$

with boundary conditions

$$v(t, R_i(t)) = R_i(t) - R_i(t_0); (i = 1, 2),$$

$$p_i(t) + (\lambda + 2\mu) \frac{\partial v}{\partial r} + \lambda R_i^{-1} v = 0,$$

where  $r = R_i(t); (i = 1, 2)$  and initial data

$$v|_{t=t_0} = v_0(r); \quad v_t|_{t=t_0} = v_1(r); \quad R_i|_{t=t_0} = R_i(t_0).$$

We use given values

$$p_i(t), \quad p_i(t_0) = p_{i0}, \quad R_i(t_0) = R_{i0}, \quad \lambda, \mu, p_0$$

and coordinated conditions  $v_0(R_{i0}) = 0, (i = 1, 2)$ .

According to the above group classification, the equation (1) admits the operators:

$$\begin{aligned} X_1 &= \partial_t; \\ X_2 &= t\partial_t + r\partial_r - \frac{1}{2}v\partial_v; \\ X_3 &= \frac{1}{2}(t^2 + r^2)\partial_t + r\partial_r - \frac{1}{2}tv\partial_v; \\ X_4 &= v\partial_v; \\ Y_g &= g(t, r)\partial_v; \quad g_u = g_{rr} + r^{-1}g_r - r^{-2}g \end{aligned} \quad (2)$$

## 2. Problem Setting

Here we consider invariants solutions with respect to three different two-dimensional sub algebras of the algebra (2):

- A.  $X_2 + \alpha X_4$ ;
- B.  $X_3 + \alpha X_4$
- C.  $X_1 + X_3 + \alpha X_4$ ;

### A. $X_2 + \alpha X_4$ Invariant Solutions Class

Invariant solution has the form

$$v = rt^{v+1}y(x), \quad x = r^2t^{-2}, \quad v = \alpha - \frac{5}{2}$$

with special boundary

$$\begin{aligned} p_j(t) &= -2(\mu + \lambda)t^{\alpha-\frac{5}{2}}y(R_j^2t^{-2}) - \\ & - (\lambda + 2\mu)R_j^2t^{\alpha-\frac{5}{2}}y'(R_j^2t^{-2}), \quad j = 1, 2, \end{aligned}$$

where  $R_j t^{\alpha-\frac{3}{2}}y(R_j^2t^{-2}) = R_j - R_{j0}$

and special initial data

$$\begin{aligned} v_0(r) &= rt_0^{\alpha-\frac{5}{2}}y(r^2t_0^{-2}), \\ v_1(r) &= (\alpha - \frac{3}{2})rt_0^{\alpha-\frac{5}{2}}y(r^2t_0^{-2}) - 2r^3t_0^{\alpha-\frac{5}{2}}y'(r^2t_0^{-2}). \end{aligned}$$

We obtain the problem on characteristic value by the equation

$$x(1-x)y'' + \left( (v - \frac{1}{2})x + 2 \right) y' - \frac{1}{4}v(v+1)y = 0.$$

Here  $y(a) = y(b) = 0$ , so that  $a = R_{10}^2t_0^{-2}$ ,  $b = R_{20}^2t_0^{-2}$ .

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### B. $X_3 + \alpha X_4$ Invariant solutions class

Invariant solution has the form

$$v = r^{-\frac{1}{2}} e^{\frac{-2a}{t+r}} \varphi(q),$$

where  $q = r - t^2 r^{-1}$ , and is determined by the equation

$$q^2 \varphi'' + 2(q + 2\alpha) \varphi' - \frac{3}{4} \varphi = 0, \quad \varphi(a) = \varphi(b) = 0,$$

so that  $a = R_{10} - \frac{t_0^2}{R_{10}^0}$ ,  $b = R_{20} - \frac{t_0^2}{R_{20}^0}$

Using the change  $\varphi(q) = \psi(p) p^b e^{-2\alpha p}$  with  $p = -q^{-1}$ ,  $b(b-1) = \frac{3}{4}$  we obtain the boundary-value problem on characteristic value  $\alpha^2$ .

$$p\psi'' + 2b\psi' - 4\alpha^2 p\psi = 0,$$

$$\psi\left(\frac{R_{10}}{R_{10}^2 - t_0^2}\right) = \psi\left(\frac{R_{20}}{R_{20}^2 - t_0^2}\right) = 0$$

The solution's equation is  $\psi = pZ_\nu(2\alpha ip)$ , where  $Z_\nu = C_1 J_\nu + C_2 Y_\nu$  - general solution of Bessel equation,  $C_1, C_2$  - arbitrary constants.

### C. $X_1 + X_3 + \alpha X_4$ Invariant Solutions Class

Invariant solution has the form

$$v = r^{-\frac{1}{2}} e^{\frac{\sqrt{2}\alpha \arctg \frac{t+r}{\sqrt{2}}}{\sqrt{2}}} \varphi(p),$$

where  $p = \frac{t^2 + 2}{r} - r$ .

Obtained equation  $(p^2 + 8)\varphi'' + 2(p - 2\alpha)\varphi' - \frac{3}{4}\varphi = 0$  with change  $\varphi = \eta\left(\frac{1}{2} - \frac{i}{4\sqrt{2}}p\right)$ , was led to hypergeometric equation

$$\xi(\xi - 1)\eta'' + \left(2\xi - \left(\frac{i}{\sqrt{2}} - 1\right)\right)\eta' - \frac{3}{4}\eta = 0,$$

with boundary conditions  $\eta(a_i) = \eta(b_i) = 0$ ,

where

$$a_i = \frac{1}{2} - i \frac{1}{4\sqrt{2}} \left( \frac{t_0^2 + R_{10}^2}{R_{10}} - R_{10} \right),$$

$$b_i = \frac{1}{2} - i \frac{1}{4\sqrt{2}} \left( \frac{t_0^2 + R_{20}^2}{R_{20}} - R_{20} \right)$$

### References

1. Abdrakhmanova R.P., Khabirov C.V. "Displacements in the elastic cylinder under pressure". The Siberian industrial mathematics magazine, Novosibirsk, Russia, 2000.